# A Comparison of Carbon Monoxide and Nitrogen as Ligands in Transition Metal Complexes 

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#### Abstract

The bonding characteristics of CO and $\mathbf{N}_{2}$ in transition metal complexes are investigated on the basis of comparative molecular orbital calculations for $\operatorname{Cr}(\mathrm{CO})_{6}$ and $\operatorname{Cr}\left(\mathrm{N}_{2}\right)_{6}$. The results suggest that the differences in $\pi$ acceptor ability of the $\mathrm{N}_{2}$ ligand compared to CO are consequences of the off-diagonal matrix element between the metal d orbitals and the $\pi$ antibonding orbital of the ligand moieties. The $\sigma$ bonding interactions can be mainly characterized as electron donation to the metal from orbitals substantially localized on the atoms adjacent to the metal. However, these orbitals have sufficient $\sigma$ antibonding properties with respect to the ligand $\sigma$ bond character that some degree of "within ligand" $\sigma$ bond strengthening on complex formation cannot be ruled out. The results suggest that some of the apparent conflicts in $\pi$ acceptor ability of the two ligands may be attributed to combined $\sigma-\pi$ factors.


TThere has been a recent surge of interest in the synthesis and study of transition metal complexes of molecular nitrogen. ${ }^{2-7}$ Nitrogen was long thought to have no affinity for transition metals on the basis of the inertness of gaseous nitrogen itself, but it now seems that it was merely necessary to develop a sufficiently clever synthetic route to nitrogen complexes.

Along with the progress in synthetic routes to these complexes there has developed some disagreement as to the $\pi$ acceptor ability of $\mathrm{N}_{2}$ relative to CO. Thus, Collman, et al., ${ }^{5}$ claim on the basis of relative changes in $\mathrm{N}_{2}$ and CO stretching frequencies in analogous compounds that " $\mathrm{N}_{2}$ is a more powerful $\pi$ acid than CO." In the same paper they also state that "Nitrogen is similar to $\mathrm{NO}^{+}$inasmuch as both are strong $\pi$ acids and weak $\sigma$ donors." Conversely, Bancroft, et al., ${ }^{6}$ state as a consequence of Mössbauer studies that " CO is an appreciably better $\sigma$ donor and/or $\pi$ acceptor than $\mathrm{N}_{2}$."

From a theoretical standpoint, CO and $\mathrm{N}_{2}$ are interesting ligand species. As free molecules, the calculated orbital energies of the two species ${ }^{8}$ are surprisingly similar if one considers that the transfer of a proton from one nucleus to the other is involved in comparing the two isoelectronic molecules. Since $\mathrm{N}_{2}$ forms complexes which are structurally similar ${ }^{9,10}$ to those of CO , a comparison of the electronic structures of analogous complexes of CO and $\mathrm{N}_{2}$ would be enlightening. Unfortunately the other ligands contained in known analogous species are of such complexity and the species are of such low molecular symmetry that theoretical computations are impossible without making severe simplifications in the calculations. Consequently, it was decided to undertake comparative calculations on the well-known compound, $\operatorname{Cr}(\mathrm{CO})_{6}$, and the hypothetical
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molecule, $\operatorname{Cr}\left(\mathrm{N}_{2}\right)_{6}$. Not only does the high symmetry of this latter species make the calculations more tractable, but it permits one to focus attention on the bonding properties of $\mathrm{N}_{2}$ relative to CO without the complexities introduced by the presence of other ligand species.

## Calculational Method

Since the publication of the earlier results from this laboratory ${ }^{11}$ for $\mathrm{Cr}(\mathrm{CO})_{6}$, minor changes have been made in the computational method in order to simplify the calculational procedure. For example, recent investigations ${ }^{12}$ have indicated that the three-center nuclear attraction integral can be very well approximated by

$$
\begin{equation*}
\left(\phi_{\mathrm{a}}\left|\frac{q_{\mathrm{v}}}{r_{\mathrm{v}}}\right| \chi_{\mathrm{b}}\right) \approx q_{\mathrm{v}} \frac{S\left(\phi^{\mathrm{a}}, \chi^{\mathrm{b}}\right)}{2}\left[\frac{1}{R_{\mathrm{av}}}+\frac{1}{R_{\mathrm{bv}}}\right] \tag{1}
\end{equation*}
$$

where $S\left(\phi^{\mathrm{a}}, \chi^{\mathrm{b}}\right)$ is the overlap integral of the functions on centers a and $\mathrm{b} ; q_{\mathrm{v}}$ is the charge on center $\mathrm{v}, R_{\mathrm{av}}$ and $R_{\mathrm{bv}}$ are the internuclear distances between the a and v centers and $b$ and $v$ centers, respectively.

The basis functions and internuclear distances for $\mathrm{Cr}(\mathrm{CO})_{6}$ were the same as those used previously. ${ }^{11}$ For purposes of comparison, the eigenvalues of the occupied orbitals obtained previously and in the present work are listed in Table I. The interpretations presented in our earlier communication ${ }^{11}$ are completely unaffected by the small deviations in the eigenvalues listed in the table. As would be expected in the case of such close agreement, the eigenfunctions of the two sets of calculations are very similar as well.

In accord with the X-ray structural analysis ${ }^{10}$ of Ru$\left(\mathrm{NH}_{3}\right)_{5} \mathrm{~N}_{2}{ }^{2+}$ which shows that the $\mathrm{N}_{2}$ ligand bonds end-on through only one nitrogen, the calculations on the hypothetical $\mathrm{Cr}\left(\mathrm{N}_{2}\right)_{6}$ assumed a structure analogous to $\mathrm{Cr}(\mathrm{CO})_{6}$. The chromium-carbon distance was set at the known distance of $1.92 \AA .^{13}$ The chromiumnitrogen distance was also taken as $1.92 \AA$ so that a comparison of CO and $\mathrm{N}_{2}$ as ligands would be unfettered by bond length changes. This invariance of $\mathrm{Cr}-\mathrm{N}$ and $\mathrm{Cr}-\mathrm{C}$ bond lengths is not incompatible with

[^0]Table I. Eigenvalues of Occupied Orbitals of $\mathrm{Cr}(\mathrm{CO})_{6}$ and $\operatorname{Cr}\left(\mathrm{N}_{2}\right)_{6}{ }^{a}$

|  | $\mathrm{Cr}(\mathrm{CO})_{6}$, <br> previous <br> work $^{b}$ | $\mathrm{Cr}(\mathrm{CO})_{6}$, <br> present <br> work | $\mathrm{Cr}\left(\mathrm{N}_{2}\right)_{6}$, <br> present <br> work |
| :---: | :---: | :---: | :---: |
| Orbital | -37.22 | -38.00 | -37.33 |
| $2 \mathrm{a}_{1 \mathrm{~g}}$ | -18.51 | -18.78 | -21.31 |
| $2 \mathrm{a}_{1 \mathrm{~g}}$ | -15.23 | -17.28 | -17.46 |
| $3 \mathrm{a}_{\mathrm{gg}}$ | -37.33 | -38.79 | -38.23 |
| $1 \mathrm{e}_{\mathrm{g}}$ | -19.01 | -19.37 | -23.14 |
| $2 \mathrm{e}_{\mathrm{g}}$ | -16.14 | -17.42 | -17.51 |
| $3 \mathrm{e}_{\mathrm{g}}$ | -37.21 | -38.14 | -37.39 |
| $1 \mathrm{t}_{1 \mathrm{u}}$ | -17.69 | -17.35 | -20.84 |
| $2 \mathrm{t}_{\mathrm{uv}}$ | -14.97 | -16.87 | -17.47 |
| $3 \mathrm{t}_{\mathrm{lu}}$ | -14.02 | -15.24 | -16.75 |
| $4 \mathrm{t}_{1 \mathrm{u}}$ | -14.95 | -15.39 | -16.83 |
| $1 \mathrm{t}_{\mathrm{tu}}$ | -14.90 | -15.36 | -16.86 |
| $1 \mathrm{t}_{1 \mathrm{~g}}$ | -16.24 | -15.80 | -17.10 |
| $1 \mathrm{t}_{2 \mathrm{~g}}$ | -8.19 | -8.30 | -7.60 |
| $2 \mathrm{t}_{2 \mathrm{~g}}$ |  |  |  |

${ }^{a} \mathrm{In} \mathrm{eV} . \quad{ }^{b}$ Reference 11.
orbital basis set used to carry out the calculations. The presentation of the eigenvectors in terms of ligand MO participation will aid in the discussion of the bonding characteristics of the species.

Even a cursory examination of the eigenvalues of $\mathrm{Cr}(\mathrm{CO})_{6}$ and $\mathrm{Cr}\left(\mathrm{N}_{2}\right)_{6}$ given in Table I suggests that the stability of the former cannot be deduced on the basis of orbital energies. While the stability of a compound requires the consideration of many factors, such as the stability of the separated moieties, it is nevertheless somewhat surprising that the eigenvalues in the two cases are so similar.

It is informative to examine the stabilization of the metal $3 \mathrm{~d} \pi$ orbitals as a consequence of interaction with the ligand $1 \pi$ and $2 \pi$ orbitals. Table IV summarizes the pertinent information. Notice that the quantity, $E\left(2 \mathrm{t}_{2 \mathrm{~g}}\right)-F(3 \mathrm{~d} \pi, 3 \mathrm{~d} \pi)$, for $\mathrm{Cr}(\mathrm{CO})_{6}$ is -2.49 eV , while its value for $\operatorname{Cr}\left(\mathrm{N}_{2}\right)_{6}$ is only -1.18 eV . Whether or not the

Table II. Eigenvectors and Eigenvalues of Occupied Orbitals of CO and $\mathrm{N}_{2}{ }^{a}$

| CO | 2 sC | 2p $\sigma \mathrm{C}$ | 2sO | $2 \mathrm{p} \sigma \mathrm{O}$ | $2 \mathrm{p} \pi \mathrm{C}$ | $2 \mathrm{p} \pi \mathrm{O}$ | Eigenvalue ${ }^{\text {b }}$ | SCF eigenvalue ${ }^{\text {c }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $3 \sigma^{d}$ | 0.285 | 0.237 | 0.698 | 0.101 |  |  | -38.60 | -40.78 |
| $4 \sigma$ | 0.223 | 0.169 | -0.485 | 0.813 |  |  | -17.95 | -19.93 |
| $5 \sigma$ | 0.785 | -0.617 | -0.070 | -0.144 |  |  | -14.63 | -13.08 |
| $1 \pi$ |  |  |  |  | 0.494 | 0.752 | -16.24 | -15.86 |
| $2 \pi^{\text {e }}$ |  |  |  |  | 0.909 | -0.711 | -1.32 | 7.09 |
| $\mathrm{N}_{2}$ | $2 \mathrm{sN} \mathrm{N}_{1}$ | $2 \mathrm{p} \sigma \mathrm{N}_{1}$ | $2 \mathrm{sN}_{2}$ | $2 \mathrm{p} \sigma \mathrm{N}_{2}$ | $2 \mathrm{p} \pi \mathrm{N}_{1}$ | $2 \mathrm{p} \pi \mathrm{N}_{2}$ | Eigenvalue ${ }^{\text {b }}$ | SCF eigenvalue ${ }^{\text {c }}$ |
| $2 \sigma_{\mathrm{g}}{ }^{\text {d }}$ | 0.495 | 0.206 | 0.495 | 0.206 |  |  | -36.30 | - 39.52 |
| $2 \sigma_{\mathrm{u}}$ | 0.589 | -0.398 | -0.589 | 0.398 |  |  | -17.42 | -19.88 |
| $3 \sigma_{\mathrm{g}}$ | 0.370 | -0.614 | 0.370 | -0.614 |  |  | -16.46 | -14.82 |
| $1 \pi_{u}$ |  |  |  |  | 0.624 | 0.624 | -16.09 | -15.77 |
| $1 \pi_{\mathrm{g}}{ }^{\text {e }}$ |  |  |  |  | 0.835 | -0.835 | 1.71 | 7.43 |

${ }^{a}$ The coordinate systems are such that $\sigma$ overlap integrals are positive. ${ }^{b}$ All eigenvalues are listed in units of eV. ${ }^{\circ}$ SCF results are taken from ref 8. ${ }^{d}$ Our method treated the 1 s functions as part of the core. The orbital designation used here is in accord with the SCF results which treat the 1 s functions as part of the basis set. ${ }^{e}$ The $2 \pi$ orbitals are unoccupied but are included here because of their significance in the bonding to the metal atom.
the X-ray diffraction results of Davis, Payne, and Ibers ${ }^{14}$ on $\operatorname{Co}\left(\mathrm{N}_{2}\right) \mathrm{H}\left[\mathrm{P}\left(\mathrm{C}_{6} \mathrm{H}_{5}\right)_{3}\right]_{3}$, which showed that the observed $\mathrm{Co}-\mathrm{N}_{2}$ distance is similar to $\mathrm{Co}-\mathrm{CO}$ distances in various cobalt carbonyl complexes. The nitrogen AO's and the nitrogen-nitrogen distance were the same as those used by Ransil ${ }^{8}$ in his SCF calculation on $\mathrm{N}_{2}$ using the Slater basis set.

## Results and Discussion

Because they will be convenient for later discussion, the eigenvectors and eigenvalues of the free ligand species, CO and $\mathrm{N}_{2}$, are given in Table II. These results are in reasonable accord with rigorous SCF values, ${ }^{8}$ particularly when one considers the approximations involved in our calculational method. The results affirm the previously stated similarity in eigenvalues for the two species.

The eigenvalues for $\mathrm{Cr}(\mathrm{CO})_{6}$ and $\mathrm{Cr}\left(\mathrm{N}_{2}\right)_{6}$ are presented in Table I. Since our previous work ${ }^{11}$ tabulated the eigenvectors of $\mathrm{Cr}(\mathrm{CO})_{6}$, which are essentially equivalent to those obtained in this work, only the $\operatorname{Cr}\left(\mathrm{N}_{2}\right)_{6}$ eigenvectors are tabulated in Table III. As with $\mathrm{Cr}-$ $(\mathrm{CO})_{6}$, the eigenvectors of $\mathrm{Cr}\left(\mathrm{N}_{2}\right)_{6}$ are reported in terms of the free ligand basis functions, i.e., the functions listed in Table II, rather than in terms of the atomic

[^1]$2 t_{2 g}$ level is higher or lower in energy than the diagonal matrix element, $F(3 \mathrm{~d} \pi, 3 \mathrm{~d} \pi)$, depends upon the competition between the $F(3 \mathrm{~d} \pi, 1 \pi)$ and $F(3 \mathrm{~d} \pi, 2 \pi)$ interactions. In $\mathrm{Cr}(\mathrm{CO})_{6}$, these are -4.34 and -7.36 eV , respectively. However, the corresponding terms in $\mathrm{Cr}\left(\mathrm{N}_{2}\right)_{6}$ are -4.76 and -5.77 eV .

The reason for the reduction of $F(3 \mathrm{~d} \pi, 2 \pi)$ in the case of the nitrogen complex is apparent upon consideration of the form of the $2 \pi$ ligand wavefunction and the resultant matrix element. A general form of the function is

$$
\begin{equation*}
\psi(2 \pi)=c_{1} \phi_{1}-c_{2} \phi_{2} \tag{2}
\end{equation*}
$$

in which the negative sign is consequence of the antibonding character of the function (see Table II); $\phi_{1}$ is a normalized symmetry adapted linear combination of wave functions on those atoms adjacent to the chromium atom, and $\phi_{2}$ is a similar set of functions on the more distant atoms. Expressed in this way, the coefficients, $c_{1}$ and $c_{2}$, can be taken directly from Table II. The matrix element then has the form

$$
\begin{equation*}
F(3 \mathrm{~d} \pi, 2 \pi)=c_{1}\langle 3 \mathrm{~d} \pi| F\left|\phi_{1}\right\rangle-c_{2}\langle 3 \mathrm{~d} \pi| F\left|\phi_{2}\right\rangle \tag{3}
\end{equation*}
$$

Because of the greater proximity of atom one to the metal atom, $\left\langle 3 \mathrm{~d} \pi \mid F \dot{\phi}_{1}\right\rangle$ will be larger than $\langle 3 \mathrm{~d} \pi| F\left|\phi_{2}\right\rangle$ in both ligands. In free $\mathrm{N}_{2}$, the two coefficients, $c_{1}$ and $c_{2}$, are identical and the total matrix element re-

Table III. Eigenvectors of Occupied Orbitals of $\mathrm{Cr}\left(\mathrm{N}_{2}\right)_{6}$ in a Basis of Free $\mathrm{N}_{2}$ MO's

|  | $3 \sigma$ | $4 \sigma$ | $5 \sigma$ | $6 \sigma$ | 4s |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $1 a_{1 g}$ | 1.0097 | 0.0234 | 0.0102 | 0.0015 | -0.0552 |  |  |
| $2 a_{1 g}$ | $-0.0821$ | 0.5828 | 0.4585 | 0.0201 | 0.2993 |  |  |
| $3 a_{1 g}$ | -0.0042 | -0.6808 | 0.7539 | 0.0054 | 0.0535 |  |  |
|  | $3 \sigma$ | $4 \sigma$ | $5 \sigma$ | $6 \sigma$ | $3 \mathrm{~d} \sigma$ | $4 \mathrm{~d} \sigma$ |  |
| $1 e_{g}$ | 1.0643 | 0.0321 | 0.0001 | -0.0056 | -0.0108 | -0.1648 |  |
| $2 \mathrm{e}_{\mathrm{g}}$ | -0.0613 | 0.5117 | 0.4962 | 0.0117 | 0.3855 | 0.1393 |  |
| $3 \mathrm{e}_{\mathrm{g}}$ | $-0.0005$ | -0.6961 | 0.6968 | -0.0046 | 0.0162 | 0.0338 |  |
|  | $1 \pi$ | $2 \pi$ | $3 \mathrm{~d} \pi$ | $4 \mathrm{~d} \pi$ |  |  |  |
| $1 \mathrm{t}_{2 \mathrm{~g}}$ | 0.9757 | 0.0592 | 0.1519 | -0.0634 |  |  |  |
| $2 t_{2 g}$ | $-0.3386$ | 0.4774 | 0.7809 | 0.0635 |  |  |  |
|  |  | $4 \sigma$ |  | $6 \sigma$ |  |  |  |
| $1 \mathrm{t}_{\text {u }}$ | $1.0272$ | $0.0214$ | $0.0047$ | $-0.0019$ | $-0.0095$ | $-0.0037$ | $-0.0843$ |
| $2 \mathrm{t}_{\text {lu }}$ | $-0.0660$ | 0.6032 | 0.5019 | 0.0147 | 0.0417 | -0.0075 | 0.3139 |
| $3 \mathrm{t}_{\text {lu }}$ | $-0.0050$ | -0.6720 | 0.7385 | 0.0052 | -0.0154 | $-0.0036$ | 0.0502 |
| $4 t_{\text {lu }}$ | 0.0032 | $-0.0182$ | 0.0497 | -0.0020 | 1.0038 | 0.0376 | 0.0221 |
| $1 \mathrm{t}_{2 \mathrm{u}}$ | $\begin{aligned} & 1 \pi \\ & 0.9993 \end{aligned}$ | $\begin{aligned} & 2 \pi \\ & 0.0426 \end{aligned}$ |  |  |  |  |  |
| $1 \mathrm{t}_{1 \mathrm{~g}}$ | 1.0000 | 0.0417 |  |  |  |  |  |

Table IV. $T_{2 g}$ Representation: Matrices, ${ }^{a} 2 \mathrm{t}_{2 \mathrm{~g}}$ Eigenfunctions, and Eigenvalues

|  | $1 \pi$ | $2 \pi$ | $3 \mathrm{~d} \pi$ | $4 \mathrm{~d} \pi$ |
| :---: | :---: | :---: | :---: | :---: |
| $F$ Matrix for $\mathrm{Cr}(\mathrm{CO})_{6}$ |  |  |  |  |
| $1 \pi$ | $-15.34$ | $-1.22$ | $-4.34$ | -3.05 |
| $2 \pi$ |  | $-2.34$ | $-7.36$ | -5.48 |
| $3 \mathrm{~d} \pi$ |  |  | $-5.81$ | 0.00 |
| $4 \mathrm{~d} \pi$ |  |  |  | 10.09 |
| $\begin{gathered} \psi\left(2 \mathrm{t}_{2 \mathrm{~g}}\right)=-0.34(1 \pi)+0.54(2 \pi)+0.69(3 \mathrm{~d} \pi)+0.04(4 \mathrm{~d} \pi) \\ E\left(2 \mathrm{t}_{2 \mathrm{~g}}\right)=-8.30 \mathrm{eV} \end{gathered}$ |  |  |  |  |
|  |  |  |  |  |
| $1 \pi$ | $-16.71$ | -0.99 | -4.76 | -4.09 |
| $2 \pi$ |  | -0.76 | -5.77 | -5.57 |
| $3 \mathrm{~d} \pi$ |  |  | $-6.42$ | 0.00 |
| $4 \mathrm{~d} \pi$ |  |  |  | 9.42 |

$\psi\left(2 \mathrm{t}_{2 \mathrm{~g}}\right)=-0.34(1 \pi)+0.48(2 \pi)+0.78(3 \mathrm{~d} \pi)+0.06(4 \mathrm{~d} \pi)$
$E\left(2 \mathrm{t}_{2 \mathrm{~g}}\right)=-7.60 \mathrm{eV}$
${ }^{a}$ In units of eV .
flects the difference between $\langle 3 \mathrm{~d} \pi| F\left|\phi_{1}\right\rangle$ and $\langle 3 \mathrm{~d} \pi| F\left|\phi_{2}\right\rangle$. Since $\mathrm{N}_{2}$ and CO are isoelectronic, the difference in the two molecules lies in the transfer of a proton from atom 1 to atom 2. Hence $\phi_{1}$ on the carbon atom should be more diffuse while $\phi_{2}$ on the oxygen becomes more contracted, which should increase the magnitude of $\langle 3 \mathrm{~d} \pi| F\left|\phi_{1}\right\rangle$ and decrease $\langle 3 \mathrm{~d} \pi| F\left|\phi_{2}\right\rangle$ in CO compared to $\mathrm{N}_{2}$. Furthermore, in CO the $c_{1}$ coefficient is larger than $c_{2}, 0.91$ compared to 0.84 . Both effects increase the first and decrease the second term in eq 3 , which results in a substantial increase in the value of the total matrix element for CO. These qualitative arguments have also been considered by Jaffé and Orchin ${ }^{15}$ and are now confirmed by the values given in Table V . The decreased $3 \mathrm{~d} \pi-2 \pi$ interaction coupled with the greater separation of the $3 \mathrm{~d} \pi$ and $2 \pi$ diagonal matrix elements accounts for the decreased $2 \pi$ and increased $3 \mathrm{~d} \pi$ character in the $2 \mathrm{t}_{2 \mathrm{~g}}$ molecular orbital for the $\mathrm{Cr}-$ $\left(\mathrm{N}_{2}\right)_{6}$ species.

In the results ${ }^{11}$ on the isoelectronic series, $\mathrm{V}(\mathrm{CO})_{6}{ }^{-}$, $\mathrm{Cr}(\mathrm{CO})_{6}$, and $\mathrm{Mn}(\mathrm{CO})_{6}{ }^{+}$, it was noted that the decreasing $2 \pi$ participation in the $2 \mathrm{t}_{2 \mathrm{~g}}$ molecular orbital was primarily a consequence of an increasing separation

[^2]Table V. Contributions to $F(3 \mathrm{~d} \pi, 2 \pi)$ in the $\mathrm{T}_{28}$ Representation

|  | CO | $\mathrm{N}_{2}$ |
| :--- | :--- | :--- |
| $c_{1}$ | 0.909 | 0.835 |
| $c_{2}$ | $0.711^{a}$ | $0.835^{a}$ |
| $\langle 3 \mathrm{~d} \pi\| F\left\|\phi_{1}\right\rangle, \mathrm{eV}$ | -8.27 | -7.20 |
| $\langle 3 \mathrm{~d} \pi\| F\left\|\phi_{2}\right\rangle, \mathrm{eV}$ | -0.22 | -0.31 |
| $F(3 \mathrm{~d} \pi, 2 \pi), \mathrm{eV}$ | -7.36 | -5.77 |

${ }^{a}$ The sign in front of $c_{2}$ was changed from that listed in Table II to conform to eq 3.
of the diagonal matrix elements. In this work, the difference in $2 \pi$ interaction appears to be equally a function of the decreased value of the off-diagonal matrix element.

Parenthetically, it might be noted that while the method employed here results in a stabilization of the $2 t_{2 g}$ level, in accord with long held qualitative arguments, ${ }^{16,17}$ more simplified calculational techniques such as the well-known SCCC method of Gray and coworkers, ${ }^{18}$ which employ the Wolfsberg-Helmholz method for evaluation of $F(3 \mathrm{~d} \pi, 1 \pi)$ and $F(3 \mathrm{~d} \pi, 2 \pi)$ always appear to result in a destabilization of the level. This result occurs because the off-diagonal elements are estimated by

$$
\begin{array}{r}
F(3 \mathrm{~d} \pi, 1 \pi)=K \cdot G(3 \mathrm{~d} \pi, 1 \pi)[F(3 \mathrm{~d} \pi, 3 \mathrm{~d} \pi)+ \\
F(1 \pi, 1 \pi)] / 2 \tag{4a}
\end{array}
$$

and

$$
\begin{array}{r}
F(3 \mathrm{~d} \pi, 2 \pi)=K \cdot G(3 \mathrm{~d} \pi, 2 \pi)[F(3 \mathrm{~d} \pi, 3 \mathrm{~d} \pi)+ \\
F(2 \pi, 2 \pi)] / 2 \tag{4b}
\end{array}
$$

where $G(\mathrm{i}, \mathrm{j})$ is the group overlap integral, $K$ is a factor (sometimes a function of $G(i, j)$ ) often set equal to 2.0 . The values used for $F(1 \pi, 1 \pi)$ are so much greater than those used for $F(2 \pi, 2 \pi)$ that the matrix element $F$ $(3 \mathrm{~d} \pi, 1 \pi)$ is substantially larger than $F(3 \mathrm{~d} \pi, 2 \pi)$, causing the destabilization. Such a result appears to be
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Table VI. Electron Distribution in $\operatorname{Cr}(\mathrm{CO})_{6}$ and $\mathrm{Cr}\left(\mathrm{N}_{2}\right)_{6}$

|  | $3 \mathrm{~d}\left(\mathrm{t}_{2 \mathrm{~g}}\right)$ | $3 \mathrm{~d}\left(\mathrm{e}_{\mathrm{g}}\right)$ | 4 s | 4 p | 4 d | $3 \sigma$ | $4 \sigma$ | $5 \sigma$ | $6 \sigma$ | $1 \pi$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{Cr}(\mathrm{CO})_{6}$ | 3.57 | 1.50 | 0.43 | 1.45 | -0.05 | $2.00^{a}$ | 2.00 | 1.38 | 0.00 | $4.00^{b}$ |
| $\mathrm{Cr}\left(\mathrm{N}_{2}\right)_{6}$ | 4.12 | 1.07 | 0.47 | 1.28 | -0.01 | $2.00^{a}$ | 1.74 | 1.72 | 0.00 | $4.00^{b}$ |
| For $\operatorname{Cr}\left(\mathrm{N}_{2}\right)_{6}:$ | $\theta_{4}=(1 / \sqrt{2})(4 \sigma-5 \sigma)$ population $=2.00$ |  |  |  | $0.32^{b}$ |  |  |  |  |  |
|  | $\theta_{5}=(1 / \sqrt{2})(4 \sigma+5 \sigma)$ population $=1.46$ |  |  |  |  |  |  |  |  |  |

${ }^{a}$ The ligand values are those for a single CO or $\mathrm{N}_{2}$ group. ${ }^{b}$ The maximum occupancy of a $\sigma$ orbital is 2.00 electrons. The $\pi$ orbital values given here take into account the double degeneracy of the $\pi$ interactions.
an artifact of the calculational method which has doubtful validity.

## Bonding Characteristics

Electron Distribution. As in previous work in this laboratory, the computations were carried to selfconsistency of charge and electron populations on the various atoms via a Mulliken population analysis. ${ }^{19}$ Similar calculations were carried out with the Löwdin ${ }^{20}$ method for electron distribution. While the absolute values varied in the two cases, as expected, ${ }^{21}$ the general trends in distributions between the two molecules were the same. Since much of the previous work is reported in terms of the Mulliken method, it will be used here as well.

It is worthwhile to represent the ligand charge distributions in terms of the ligand molecular orbitals rather than the individual atomic orbitals. The charge distributions are presented in Table VI. The trends are indicative on several counts.

If one were to envision a crystal field model for the metal atom configuration, it would be $\left(\mathrm{t}_{2 \mathrm{~g}}\right)^{6} \mathrm{e}_{\mathrm{g}}{ }^{0}$. The decrease in the $t_{2 g}$ occupation then becomes an estimate of the ability of the metal-ligand $\pi$ interaction to transfer charge to the ligands. Conversely, the increase of density within the $e_{g}$ metal orbitals estimates the $\sigma$ bonding donation from the ligands to the metal. (These arguments represent a somewhat simplified view which ignores the complexity introduced by the $4 p$ orbitals which are capable of both $\sigma$ and $\pi$ interaction.) The smaller $t_{2 g}$ and larger $e_{g}$ metal orbital populations for $\mathrm{Cr}(\mathrm{CO})_{6}$ clearly indicate the better $\sigma$ and $\pi$ interaction of $\mathrm{Cr}(\mathrm{CO})_{6}$ compared to $\mathrm{Cr}\left(\mathrm{N}_{2}\right)_{6}$, and display the mutual enhancement of these interactions dia the synergic effect. Notice that it is the orbital occupancy which differs much more dramatically than the total 3d populations.

The electron populations in the $1 \pi$ ligand orbitals indicate that any $\pi$ donating effect by this orbital through, for example, the $1 \mathrm{t}_{\mathrm{gg}}$ molecular orbitals is counterbalanced by back donation, so that the $1 \pi$ occupation is unchanged from that of the free ligand. The significant $\pi$ interaction, of course, involves the $2 \pi$ ligand orbitals which accept electron density via interaction with the filled metal $3 \mathrm{~d} \pi$ orbitals. The larger $2 \pi$ occupation for $\mathrm{Cr}(\mathrm{CO})_{6}$ than for $\mathrm{Cr}\left(\mathrm{N}_{2}\right)_{6}$ is in keeping with the $3 \mathrm{~d} \pi$ occupancy previously mentioned.

One interesting result appears upon examination of the ligand populations in the $4 \sigma$ and $5 \sigma$ orbitals. These two orbitals are completely filled in the free ligands. In $\mathrm{Cr}(\mathrm{CO})_{6}$, the $\sigma$ donation to the metal occurs via the $5 \sigma$ orbital, while in $\mathrm{Cr}\left(\mathrm{N}_{2}\right)_{6}$, both the $4 \sigma$ and $5 \sigma$ are involved to approximately the same extent. From Table
(19) R. S. Mulliken, J. Chem. Phys., 23, 1833 (1955).
(20) P. O. Löwdin, Phys, Rev., 18, 365 (1950).
(21) (a) E. R. Davidson, J. Chem. Phys., 46, 3320 (1967); (b) E. W. Stout, Jr., and P. Politzer, Theor. Chim. Acta, 12, 379 (1968).

II, one sees that the $5 \sigma$ orbital in CO is substantially localized on the carbon atom. In the $\mathrm{N}_{2}$ complex, the $4 \sigma$ and $5 \sigma$ orbitals (labeled $2 \sigma_{\mathrm{u}}$ and $3 \sigma_{\mathrm{g}}$ in Table II) are, by the choice of basis set, delocalized over the $\mathrm{N}_{2}$ molecule. However, if one takes linear combinations of the two functions

$$
\begin{align*}
& \theta_{4}=(1 / \sqrt{2})(4 \sigma-5 \sigma)=(1 / \sqrt{2})\left(2 \sigma_{\mathrm{u}}-3 \sigma_{\mathrm{g}}\right) \\
& \theta_{5}=(1 / \sqrt{2})(4 \sigma+5 \sigma)=(1 / \sqrt{2})\left(2 \sigma_{\mathrm{u}}+3 \sigma_{\mathrm{g}}\right) \tag{5}
\end{align*}
$$

then these two new functions are again essentially localized on nitrogen atoms 2 and 1, respectively. Their populations are also listed in Table VI and clearly indicate that, as in the carbonyl complex, the $\sigma$ donation to the metal comes essentially from the atom adjacent to the metal. The values of 1.38 for $5 \sigma$ on CO and 1.46 for $\theta_{5}$ on $\mathrm{N}_{2}$ again illustrate the weaker $\sigma$ donating ability of the $\mathrm{N}_{2}$ group in these complexes.

While the substantially localized character of the $5 \sigma$ orbital in CO and the $\theta_{5}$ orbital in $\mathrm{N}_{2}$ aid in correlating the MO results with a valence bond localized orbital interaction, it is important to keep in mind that the two orbitals are slightly antibonding with respect to the $\sigma$ framework of the two ligand species. Thus, the $5 \sigma$ orbital of free CO has a negative overlap population, -0.047 , as does $\theta_{5}$ in $\mathrm{N}_{2},-0.036$. Removal of electron density from these orbitals results in an increase in $\sigma$ bond strength for their respective molecules. Therefore, although $\pi$ bond strength in the CO complex is reduced to a greater degree than that in the $\mathrm{N}_{2}$ complex, there is a greater $\sigma$ bond strengthening in the CO complex, since a greater portion of electron density was removed from the $5 \sigma$ orbital, and the orbital possesses a larger antibonding character than the $\theta_{5}$ function.

It is clear, then, that the resultant bond strength in the two ligand species compared to their strengths as free ligands is not necessarily relatable to $\pi$ interaction only. Since such characteristics as stretching frequencies will depend on total bond strength, the $\pi$ interactions cannot be isolated from potential $\sigma$ effects. It is possible that some of the controversy over relative $\pi$ acceptor abilities of these ligands has resulted from failure to consider both changes. Furthermore, the balance between $\sigma$ bond strengthening and $\pi$ bond weakening suggests that the final bond character in mixed-ligand complexes could be dependent upon the nature of the other attached ligands. It is the suggestion of the authors that the properties of CO and $\mathrm{N}_{2}$ in $\mathrm{ML}_{5} \mathrm{X}$ (where $\mathrm{X}=\mathrm{CO}$ or $\mathrm{N}_{2}$ ) will depend upon the ligand L . Future work in this laboratory will include efforts to verify this conjecture.

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